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# Temperature dependence of photoluminescence from Mg-doped $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$ grown by liquid-phase epitaxy

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The temperature dependence of photoluminescence from the Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers on (100) GaAs substrates grown by liquid-phase epitaxy has been studied. At low temperature, the spectra show only two major emission peaks involving intrinsic recombination and conduction-band-to-acceptor transition. The intrinsic recombination dominates in the doping concentration range studied ( $1.0 \times 10^{17}$ – $7.0 \times 10^{18} \text{ cm}^{-3}$ ) above 60 K. Below 50 K, these two peaks merged with each other when the doping concentration is higher than  $1 \times 10^{18} \text{ cm}^{-3}$ . The temperature dependence of band gap in  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers determined from the photoluminescence peak energy varies as  $1.976 - [7.5 \times 10^{-4} T^2 / (T + 500)] \text{ eV}$ . For the moderately doped concentration ( $p < 1.4 \times 10^{18} \text{ cm}^{-3}$ ), the Mg acceptor ionization energy obtained from 50-K photoluminescent spectra is in the range from 37 to 40 meV.

## I. INTRODUCTION

$\text{In}_{1-x}\text{Ga}_x\text{P}$  ternary compound offers the possibility of direct recombination luminescence up to photon energy of approximately 2.2 eV at 300 K and was predicted to be a highly efficient luminescence material,<sup>1</sup> providing an opportunity for the fabrication of visible light-emitting devices. Especially, with  $x \sim 0.5$ , it can be lattice matched to GaAs, and thus has attracted much attention for the fabrication of red light-emitting devices.<sup>2-5</sup>

Mg and Zn are well known *p*-type dopants in GaAs and AlGaAs layers grown by liquid-phase epitaxy (LPE). Mg is expected to be more attractive than Zn as a *p*-type dopant in GaAs and related III-V compounds because the diffusion coefficient of Mg is about  $10^3$  times lower than that of Zn in GaAs.<sup>6,7</sup> Zn has been shown to migrate in AlGaAs at low temperature and produce compositional disorder at the AlGaAs-GaAs interface.<sup>8</sup> Recently, remarkable improvement in Mg doping efficiency, compared with that for Zn, has been observed for  $(\text{Al}_x\text{Ga}_{1-x})_{0.5}\text{In}_{0.5}\text{P}$ , grown by metalorganic chemical vapor deposition (MOCVD), especially for a high aluminum composition layers.<sup>9</sup> Suzuki *et al.* reported the low-temperature photoluminescence (PL) properties of Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  grown by MOCVD with different hole concentrations.<sup>10</sup> Chang *et al.* reported the doping properties of Mg in  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers grown by LPE.<sup>11</sup> But there is no detailed report on the temperature dependence of luminescent properties of Mg-doped InGaP material or devices.

## II. EXPERIMENT

This article reports the temperature dependence of photoluminescence of Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers with various doping concentrations.

The Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  epitaxial layers used for this study were grown on undoped semi-insulating GaAs substrates with a horizontal sliding LPE system. The growth temperature, the initial supercooling, and the cooling rate

were 785 °C, 12 °C, and 0.3 °C/min, respectively. The thickness of the layers during a growth period of 17 min was 8–10  $\mu\text{m}$ . Details of the LPE growth was described in Ref. 5.

The PL measurements were made using argon laser excitation (4880-Å line) with an average power density of  $\sim 5 \text{ W/cm}^2$  and a spot size of  $\sim 1 \text{ mm}$  in diameter. The emission spectra were analyzed by a 1-m spectrometer and detected with a silicon pin photodetector. The samples were mounted on a holder inside a cryostat and the temperature was varied and controlled by a thermalfilm heater wound near the holder and detected by a Pt temperature sensor.

## III. RESULTS AND DISCUSSIONS

Table I illustrates some of the electrical and optical properties at room temperature of the undoped and Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  samples used for this investigation. An undoped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  sample with background concentration of  $4 \times 10^{16} \text{ cm}^{-3}$  was used as a reference for the PL measurements. The fact that the room-temperature PL full width at half maximum (FWHM) of the four Mg-doped samples broadens as the carrier concentration increased is due to the evolution of the impurity levels into a band of states.

The PL spectra at 50 K of all the samples illustrated in Table I are shown in Fig. 1. These spectra are normalized to the same main peak intensity. Except for the undoped and most heavily doped samples No. 1 and No. 5, respectively, as shown in Figs. 1(a) and 1(e), all other samples exhibit three peaks denoted as A, B, and C. Peak A with the narrowest FWHM and the highest photon energy as shown in Fig. 1(a) of the undoped sample is corresponding to band-to-band (intrinsic) transition. The lower-energy peak B, the dominant transition process at temperature lower than 60 K, is due to the transition of band to acceptor level (BA). As shown in Fig. 1(b) with a hole concentration of  $1.0 \times 10^{17} \text{ cm}^{-3}$ , its position is 38 meV below

TABLE I. 300-K electrical and optical properties of the Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  samples used in this study.

Sample No.	$n$ or $p$ ( $\times 10^{17} \text{ cm}^{-3}$ )	Mobility ( $\text{cm}^2/\text{V s}$ )	PL peak energy (eV)	PL FWHM (meV)	PL intensity (AU)
1	undoped 0.4( $n$ )	1300	1.890	36.5	1.8
2	1.0( $p$ )	65	1.889	41.2	2.1
3	5.1( $p$ )	52	1.890	45.0	4.5
4	14 ( $p$ )	41	1.880	59.5	3.8
5	70 ( $p$ )	29	1.885	75.5	0.6

peak A and moves closer to peak A with increasing hole concentration. Since the relative intensity of peak B to peak A increases with hole concentration, it confirms that peak B is definitely associated with the Mg acceptor level. The weak peak C located at  $\sim 50$  meV below peak B is believed to be the phonon replica of peak B. As the hole concen-

tration is higher than  $\sim 1 \times 10^{18} \text{ cm}^{-3}$ , these two peaks A and B merged with each other and is not distinguished even at 16 K. This is due to the evolution of acceptor impurity levels into a band and merging with the valence band.

From this BA transition of peak B, the ionization energy of the Mg acceptor,  $E_A$ , can be estimated from the photon energy of peak B,  $h\nu_B$ , and  $E_g$  using Eagles equation:<sup>12</sup>

$$E_A = E_g - h\nu_B + \frac{1}{2}kT,$$

where  $E_g$  is the band-gap energy of  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  and  $h\nu_B$  is the peak energy of the Mg acceptor. The thermal-energy correction term  $\frac{1}{2}kT$  adds about 2 meV to the term of  $(E_g - h\nu_B)$  at 50 K. The ionization energy of Mg acceptor in the lightly doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  is thus calculated as 37–40 meV at 50 K.

Figure 2(a) presents the temperature dependence ( $T=16$ –300 K) of the emission peak of sample No. 2 with a hole concentration of  $1 \times 10^{17} \text{ cm}^{-3}$ . Peak B is the dominant recombination process at temperature below 60 K. As the temperature is increased, the intensity of this peak gradually decreased while that of peak A gradually increased. Above 130 K, the emission of peak B becomes negligibly weak. The decreasing of relative intensity of peak B with increasing temperature is due to the thermal release of holes from acceptor sites.

The PL spectra of sample No. 4 with a hole concentration of  $1.4 \times 10^{18} \text{ cm}^{-3}$  is shown in Fig. 2(b). At 16 K, only peaks B and C are observed and it can be explained as the increase of transition probability due to higher hole population in the acceptor level. It is relevant to note in Table I that the PL intensity at 300 K reaches its maximum in the range from  $5 \times 10^{17}$  to  $1.4 \times 10^{18} \text{ cm}^{-3}$  and then drops quickly as the hole concentration is further increased. There may be relatively deep centers formed at high Mg-doped concentrations to lower the radiative efficiency as in the case of heavily Zn-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers.<sup>13–15</sup>

Figure 3 shows the variation of the photon energies of the peaks A and B as a function of temperature for sample No. 2 with PL spectra shown in Fig. 2(a). The temperature dependence of the band gap in this figure can be expressed as the Varshni equation:<sup>16</sup>

$$E_g(T) = E_g(0) - \alpha T^2 / (T + \beta),$$

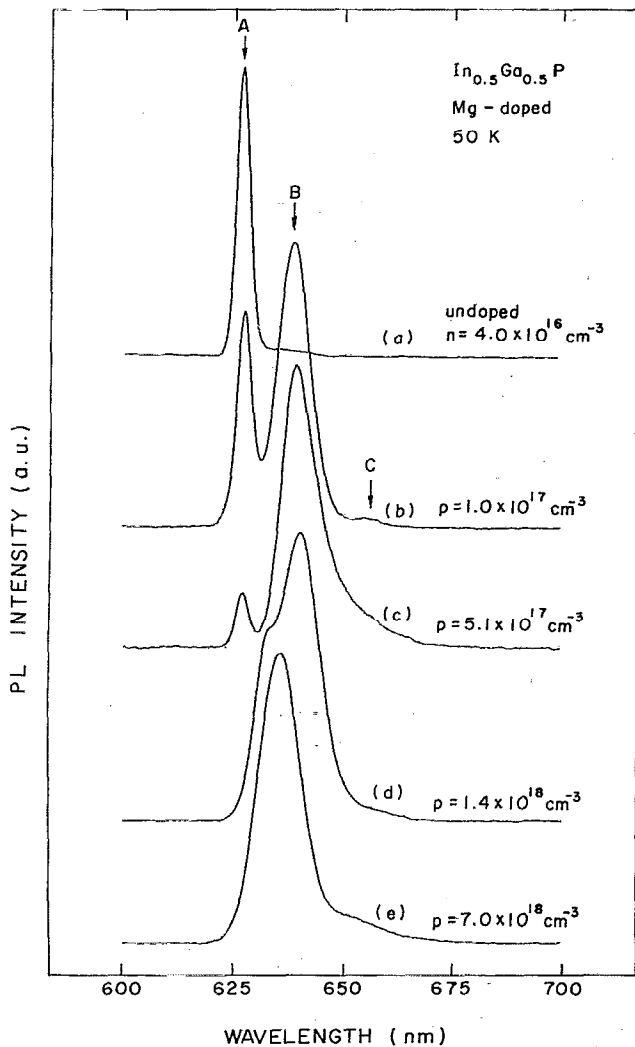


FIG. 1. 50-K PL spectra of Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers grown on GaAs with a carrier concentration of (a)  $n = 4.0 \times 10^{16} \text{ cm}^{-3}$  (undoped), (b)  $p = 1.0 \times 10^{17} \text{ cm}^{-3}$ , (c)  $p = 5.1 \times 10^{17} \text{ cm}^{-3}$ , (d)  $p = 1.4 \times 10^{18} \text{ cm}^{-3}$ , and (e)  $p = 7.0 \times 10^{18} \text{ cm}^{-3}$ . The spectra are normalized to the same main peak intensity.

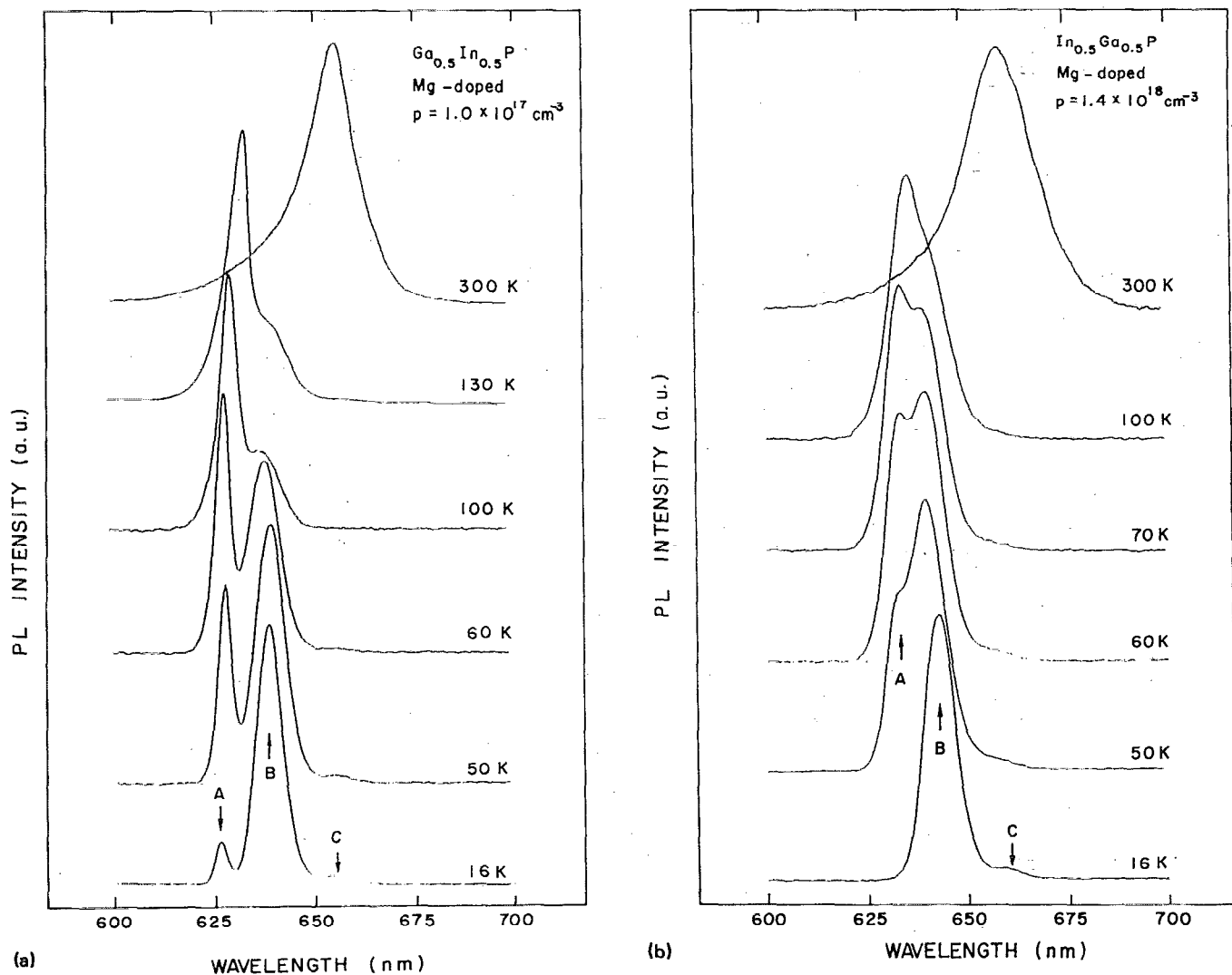


FIG. 2. Photoluminescence spectra of the Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers with (a)  $p = 1.0 \times 10^{17} \text{ cm}^{-3}$  and (b)  $p = 1.4 \times 10^{18} \text{ cm}^{-3}$  at various temperatures between 16 and 300 K showing the gradual evolution of the peaks. The spectra are normalized to the same main peak intensity.

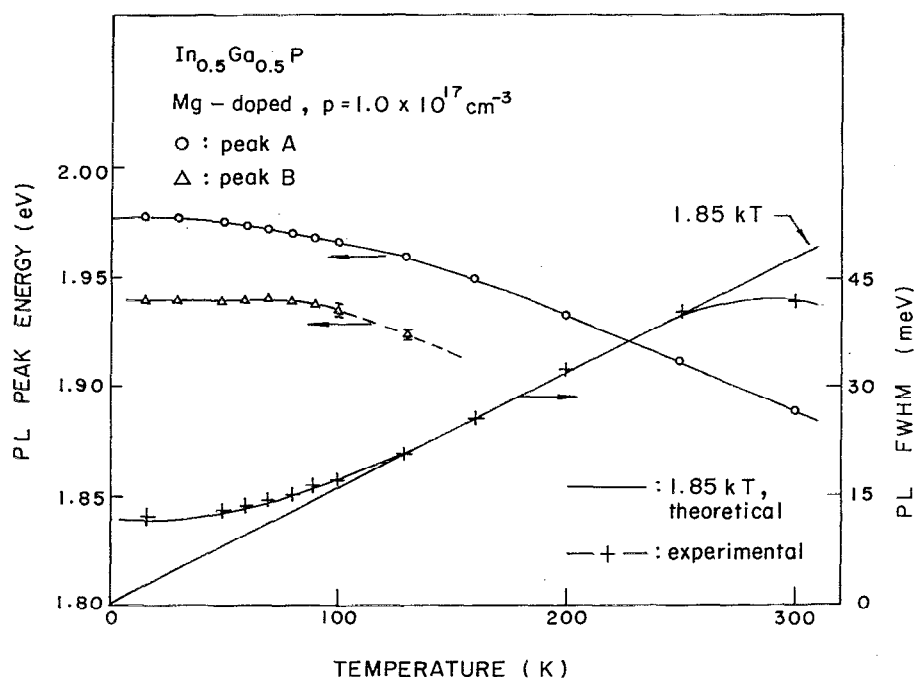


FIG. 3. Variation of the photon energies of the peaks A and B as a function of temperature for sample No. 2 with PL spectra as shown in Fig. 2 (a).

where  $E_g(0)$  is the energy gap at 0 K,  $\alpha$  and  $\beta$  are material constants. The calculated  $E_g(0)$ ,  $\alpha$ , and  $\beta$  by least-square method are 1.976 eV,  $7.5 \times 10^{-4}$  eV/K, and 500 K, respectively. Temperature dependence of the FWHM of peak A is also shown in Fig. 3. The experimental value agrees well with the theoretical value of  $1.85 kT$  (Ref. 17) as shown in Fig. 3, where  $k$  is the Boltzmann constant. The slightly smaller FWHM than the theoretical value at 300 K seems to be due to reabsorption of emitted photons with energies higher than the peak energy by the epitaxial layer,<sup>18</sup> while the somewhat larger FWHM at the temperatures lower than 100 K may be due to the existence of another radiative transition process such as (residual) donor-to-valence-band recombination very near to the band-to-band transition peak.

#### IV. CONCLUSIONS

We have demonstrated the temperature dependence of photoluminescence of Mg-doped  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  material with carrier concentrations of  $1.0 \times 10^{17}$ – $7.0 \times 10^{18} \text{ cm}^{-3}$ . The relationship between the temperature and PL relative intensities of the two major peaks termed the intrinsic transition and band-to-acceptor recombination has been investigated in detail. At temperature above 60 K, the intrinsic recombination dominates in the doping range studied. The temperature dependence of band gap in  $\text{In}_{0.5}\text{Ga}_{0.5}\text{P}$  layers determined from PL peak energy varies as  $1.976 - [7.5 \times 10^{-4} T^2 / (T + 500)] \text{ eV}$ .

#### ACKNOWLEDGMENTS

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